Magnetic reconnection for 10-moment two-fluid versus kinetic simulations E. Alec Johnson (ejohnson@math.wisc.edu) and James A. Rossmanith (rossmani@math.wisc.edu)

Abstract

We compare two-fluid simulations of the GEM challenge problem with published simulations using kinetic models (Vlasov and PIC). The tenmoment model with pressure isotropization agrees well with kinetic models in the rate of reconnection, the qualitative structure of the electron diffusion region, and the dominance of the pressure term in Ohm's law and is computationally inexpensive due to the absence of diffusive terms.

GEM challenge problem

The GEM magnetic reconnection challenge problem^a was formulated to compare the ability of different plasma models to model fast magnetic reconnection. It initiates reconnection by pinching adjacent oppositely directed field lines from their equilibrium state.

^{*a*}Birn et al, Geospace environmental modeling (GEM) magnetic reconnection challenge, Journal of Geophysical Research—Space Physics, 106:3715–3719, 2001.

Ten-moment two-fluid model

We neglected all interspecies collision terms. The ten-moment two-fluid model we used assumes (1) conservation of mass and momentum and pressure tensor evolution for each species *s*:

$$\partial_t \rho_{\rm s} + \nabla \cdot (\rho_{\rm s} \mathbf{u}_{\rm s}) = 0,$$

$$\partial_t (\rho_{\rm s} \mathbf{u}_{\rm s}) + \nabla \cdot (\rho_{\rm s} \mathbf{u}_{\rm s} \mathbf{u}_{\rm s} + \mathbb{P}_{\rm s}) = \frac{q_{\rm s}}{m_{\rm s}} \rho_{\rm s} (\mathbf{E} + \mathbf{u}_{\rm s} \times \mathbf{B}),$$

$$\partial_t \mathbb{P}_{\rm s} + \nabla \cdot (\mathbf{u}_{\rm s} \mathbb{P}_{\rm s}) + 2 \operatorname{Sym} (\mathbb{P}_{\rm s} \cdot \nabla \mathbf{u}_{\rm s}) + \nabla \cdot \mathbb{q}_{\rm s}$$

$$= 2 \operatorname{Sym} \left(\frac{q_{\rm s}}{m_{\rm s}} \mathbb{P}_{\rm s} \times \mathbf{B}\right) + \mathbb{R}_{\rm s},$$

and (2) Maxwell's equations for evolution of electromagnetic field:

> $\partial_t \mathbf{B} + \nabla \times \mathbf{E} = 0,$ $\nabla \cdot \mathbf{B} = 0,$ $\partial_t \mathbf{E} - c^2 \nabla \times \mathbf{B} = -\mathbf{J}/\epsilon,$ $\nabla \cdot \mathbf{E} = \sigma/\epsilon.$

To provide for isotropization we let $\mathbb{R}_s = \frac{1}{\tau_s} \left(\frac{1}{3} (\operatorname{tr} \mathbb{P}_s) \mathbb{I} - \mathbb{P}_s \right)$, and for the isotropization period we used $\tau_{\rm s} = \tau_0 \sqrt{\frac{\det \mathbb{P}_{\rm s}}{\rho^5}} m_{\rm s}^3$, which attempts to generalize the Braginskii closure; for the GEM problem this means that $\tau_i/\tau_e \approx (m_i/m_e)^{5/4}$. Isotropization provides hyperbolic viscosity; rapid isotropization is asymptotically equivalent to a small viscosity $\eta \approx p_{\rm s} \tau_{\rm s}$. We set $\tau_0 = 50$. We set $q_s = 0$. For conservation and shockcapturing purposes we evolve the *energy* tensor $\mathbb{E}_s := \mathbb{P}_s +$ $\rho_s \mathbf{u}_s \mathbf{u}_s$ rather than the pressure tensor. We implemented an explicit third-order discontinuous Galerkin solver in the DOG-PACK framework.



 $15.3/\Omega_i$

 $13.0/\Omega_{
m i}$

5-moment

5-moment

[LoHaShu10]

[JoRo10]

[ScGr06] H. Schmitz and R. Grauer. Kinetic Vlasov simulations of collisionless magnetic reconnection. *Physics of Plasmas*, 13, 092309 (2006).

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